

# HYPOELLIPTICITY OF THE $\bar{\partial}$ -NEUMANN PROBLEM AT EXPONENTIALLY DEGENERATE POINTS

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ABSTRACT. We prove local hypoellipticity of the complex Laplacian  $\square$  in a domain which has compactness estimates, is of finite type outside a curve transversal to the CR directions and for which the holomorphic tangential derivatives of a defining function are subelliptic multipliers in the sense of Kohn.

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## 1. INTRODUCTION

For the pseudoconvex domain  $\Omega \subset \mathbb{C}^n$  whose boundary is defined in coordinates  $z = x + iy$  of  $\mathbb{C}^n$ , by

$$(1.1) \quad 2x_n = \exp \left( -\frac{1}{(\sum_{j=1}^{n-1} |z_j|^2)^{\frac{s}{2}}} \right), \quad s > 0,$$

the tangential Kohn Laplacian  $\square_b = \bar{\partial}_b \bar{\partial}_b^* + \bar{\partial}_b^* \bar{\partial}_b$  as well as the full Laplacian  $\square = \bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}$  show very interesting features especially in comparison with the “tube domain” whose boundary is defined by

$$(1.2) \quad 2x_n = \exp \left( -\frac{1}{(\sum_{j=1}^{n-1} |x_j|^2)^{\frac{s}{2}}} \right), \quad s > 0.$$

(Here  $z_j$  have been replaced by  $x_j$  at exponent.) Energy estimates are the same for the two domains. For the problem on the boundary  $b\Omega$ , they come as

$$(1.3) \quad \|(\log \Lambda)^{\frac{1}{s}} u\|_{b\Omega} \lesssim \|\bar{\partial}_b u\|_{b\Omega}^2 + \|\bar{\partial}_b^* u\|_{b\Omega}^2 + \|u\|_{b\Omega}^2$$

for any smooth compact support form  $u \in C_c^\infty(b\Omega)^k$  of degree  $k \in [1, n-2]$ .

Here  $\log \Lambda$  is the tangential pseudodifferential operator with symbol  $\log(1 + |\xi'|^2)^{\frac{1}{2}}$ ,  $\xi' \in \mathbb{R}^{2n-1}$ , the dual real tangent space. As for the problem on the domain  $\Omega$ , one has simply to replace  $\bar{\partial}_b, \bar{\partial}_b^*$  by  $\bar{\partial}, \bar{\partial}^*$  and take norms over  $\Omega$  for forms  $u$  in  $D_{\bar{\partial}^*}$ , the domain of  $\bar{\partial}^*$ , of degree  $1 \leq k \leq n-1$ ; this can be seen, for instance, in [9]. In particular, these are superlogarithmic (resp. compactness) estimates if  $s < 1$  (resp. for

any  $s > 0$ ). A related problem is that of the local hypoellipticity of the Kohn Laplacian  $\square_b$  or, with equivalent terminology, the local regularity of the inverse (modulo harmonics) operator  $N_b = \square_b^{-1}$ . Similar is the notion of hypoellipticity of the Laplacian  $\square$  or the regularity of the inverse Neumann operator  $N = \square^{-1}$ . It has been proved by Kohn in [12] that superlogarithmic estimates suffice for local hypoellipticity of the problem both in the boundary and in the domain. (Note that hypoellipticity for the domain, [12] Theorem 8.3, is deduced from microlocal hypoellipticity for the boundary, [12] Theorem 7.1, but a direct proof is also available, [7] Theorem 5.4.) In particular, for (1.1) and (1.2), there is local hypoellipticity when  $s < 1$ .

As for the more delicate hypoellipticity, in the uncertain range of indices  $s \geq 1$ , only the tangential problem has been studied and the striking conclusion is that the behavior of (1.1) and (1.2) split. The first stays always hypoelliptic for any  $s$  (Kohn [11]) whereas the second is not for  $s \geq 1$  (Christ [4]). When one tries to relate  $(\bar{\partial}_b, \bar{\partial}_b^*)$  on  $b\Omega$  to  $(\bar{\partial}, \bar{\partial}^*)$  on  $\Omega$ , estimates go well through (Kohn [12] Section 8 and Khanh [7] Chapter 4) but not regularity. In particular, the two conclusions about tangential hypoellipticity of  $\square_b$  for (1.1) and non-hypoellipticity for (1.2) when  $s \geq 1$ , cannot be automatically transferred from  $b\Omega$  to  $\Omega$ . Now, for the non-hypoellipticity in  $\Omega$  in case of the tube (1.2) we have obtained with Baracco in [1] a result of propagation which is not equivalent but intimately related. The real lines  $x_j$  are propagators of holomorphic extendibility from  $\Omega$  across  $b\Omega$ . What we prove in the present paper is hypoellipticity in  $\Omega$  for (1.1) when  $s \geq 1$ .

**Theorem 1.1.** *Let  $\Omega$  be a pseudoconvex domain of  $\mathbb{C}^n$  in a neighborhood of  $z_o = 0$  and assume that the  $\bar{\partial}$ -Neumann problem satisfies the following properties*

- (i) *there are local compactness estimates,*
- (ii) *there are subelliptic estimates for  $(z_1, \dots, z_{n-1}) \neq 0$ ,*
- (iii)  *$\partial_{z_j} r$ ,  $j = 1, \dots, n-1$ , are subelliptic multipliers (cf. [10]).*

*Then  $\square$  is locally hypoelliptic at  $z_o$ .*

The proof follows in Section 2. It consists in relating the system on  $\Omega$  to the tangential system on  $b\Omega$  along the guidelines of [12] Section 8, and then in using the argument of [11] simplified by the additional assumption (i).

*Remark 1.2.* The domain with boundary (1.1), but not (1.2), satisfies the hypotheses of Theorem 1.1 for any  $s > 0$ : (i) is obvious, and (ii) and (iii) are the content of [11] Section 4.

Notice that  $\partial\Omega$  is given only locally in a neighborhood of  $z_o$ . We can continue  $\partial\Omega$  leaving it unchanged in a neighborhood of  $z_o$ , making it strongly pseudoconvex elsewhere, in such a way that it bounds a relatively compact domain  $\Omega \subset\subset \mathbb{C}^n$  (cf. [14]). In this situation  $\square$  is hypoelliptic at every boundary point. Also, it is well defined a  $H^0$  inverse Neumann operator  $N = \square^{-1}$ , and, by Theorem 1.1, the  $\bar{\partial}$ -Neumann solution operator  $\bar{\partial}^*N$  preserves  $C^\infty(\bar{\Omega})$ -smoothness. It even preserves the exact Sobolev class  $H^s$  according to Theorem 2.7 below. In other words, the canonical solution  $u = \bar{\partial}^*Nf$  of  $\bar{\partial}u = f$  for  $f \in \text{Ker } \bar{\partial}$  is  $H^s$  exactly at the points of  $b\Omega$  where  $f$  is  $H^s$ . The Bergman projection  $B$  also preserves  $C^\infty(\bar{\Omega})$ -smoothness on account of Kohn's formula  $B = \text{Id} - \bar{\partial}^*N\bar{\partial}$ .

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## 2. HYPOELLIPTICITY OF $\square$ AND EXACT HYPOELLIPTICITY OF $\bar{\partial}^*N$

We state properly hypoellipticity and exact hypoellipticity of a general system  $(P_j)$ .

**Definition 2.1.** (i) The system  $(P_j)$  is locally hypoelliptic at  $z_o \in b\Omega$  if

$$P_j u \in C^\infty(\bar{\Omega})_{z_o}^k \text{ for any } j \text{ implies } u \in C^\infty(\bar{\Omega})_{z_o}^k,$$

where  $C^\infty(\bar{\Omega})_{z_o}^k$  denotes the set of germs of  $k$ -forms smooth at  $z_o$ .

(ii) The system  $(P_j)$  is exactly locally hypoelliptic at  $z_o \in b\Omega$  when there is a neighborhood  $U$  of  $z_o$  such that for any pair of cut-off functions  $\zeta$  and  $\zeta'$  in  $C_c^\infty(U)$  with  $\zeta'|_{\text{supp}(\zeta)} \equiv 1$  we have for any  $s$  and for suitable  $c_s$

$$(2.1) \quad \|\zeta u\|_s^2 \leq c_s \left( \sum_j \|\zeta' P_j u\|_s^2 + \|u\|_0^2 \right), \quad u \in C^\infty(\bar{\Omega})^k \cap D_{(P_j)}.$$

If  $(P_j)$  happens to have an inverse, this is said to be locally regular and locally exactly regular in the situation of (i) and (ii) respectively.

*Remark 2.2.* By Kohn-Nirenberg [13] the assumption  $u \in C^\infty$  can be removed from (2.1). Precisely, by the elliptic regularization, one can prove that if  $\zeta' P_j u \in H^s$  and  $\zeta' u \in H^0$ , then  $\zeta u \in H^s$  and satisfies (2.1). This motivates the word “exact”, that is, Sobolev exact. Not only the local  $C^\infty$ - but also the  $H^s$ -smoothness passes from  $P_j u$  to  $u$ .

Let  $\vartheta$  be the formal adjoint of  $\bar{\partial}$  and  $\Delta = \bar{\partial}\vartheta + \vartheta\bar{\partial}$  the Laplacian; it acts on forms by the action of the usual Laplacian on its coefficients.

If  $u \in D_{\square}$ , then  $\square u = \Delta u$ . We first prove exact hypoellipticity of the system  $(\bar{\partial}, \bar{\partial}^*, \Delta)$ ; hypoellipticity of  $\square$  itself will follow by the method of Boas-Straube.

**Theorem 2.3.** *In the situation of Theorem 1.1, we have, for a neighborhood  $U$  of  $z_o$  and for any couple of cut-off  $\zeta$  and  $\zeta'$  with  $\zeta'|_{\text{supp } \zeta} \equiv 1$*

$$(2.2) \quad \|\zeta u\|_s^2 \lesssim \|\zeta' \bar{\partial} u\|_s^2 + \|\zeta' \bar{\partial}^* u\|_s^2 + \|\zeta' \Delta u\|_{s-2}^2 + \|u\|_0^2, \quad u \in D_{\bar{\partial}^*}.$$

*In particular, the system  $(\bar{\partial}, \bar{\partial}^*, \Delta)$  is exactly locally hypoelliptic at  $z_o = 0$ .*

*Remark 2.4.* The hypoellipticity of  $\square_b$  under (ii) and (iii) of Theorem 1.1 is proved by Kohn in [11]. It does not require (i) but it is not exact hypoellipticity (the neighborhood  $U$  of (2.1) depends on  $s$ ). However, inspection of his proof shows that, if (i) is added, then in fact (2.1) holds for  $(P_j) = \square_b$ . Our proof consists in a reduction to the tangential system.

*Proof.* We proceed in several steps which are highlighted in two intermediate propositions. We use the standard notation  $Q(u, u)$  for  $\|\bar{\partial} u\|_0^2 + \|\bar{\partial}^* u\|_0^2$  and some variants as, for an operator  $Op$ ,  $Q_{Op}(u, u) := \|Op \bar{\partial} u\|_0^2 + \|Op \bar{\partial}^* u\|_0^2$ ; most often, in our paper,  $Op$  is chosen as  $\Lambda^s \zeta'$ . We decompose a form  $u$  as

$$\begin{cases} u = u^\tau + u^\nu, \\ u^\tau = u^{\tau+} + u^{\tau-} + u^{\tau 0}, \end{cases}$$

where the first is the decomposition in tangential and normal component and the second is the microlocal decomposition  $u^{\tau 0} = \Psi^{\pm 0} u^\tau$  in which  $\Psi^{\pm 0}$  are the tangential pseudodifferential operators whose symbols  $\psi^{\pm 0}$  are a conic decomposition of the unity in the space dual to  $\mathbb{R}^{2n-1}$  the real orthogonal to  $\partial r$  (cf. Kohn [12]). We begin our proof by remarking that any of the forms  $u^\# = u^\nu, u^{\tau-}, u^{\tau 0}$  enjoys elliptic estimates

$$(2.3) \quad \|\zeta u^\#\|_s^2 \lesssim \|\zeta' \bar{\partial} u^\#\|_{s-1}^2 + \|\zeta' \bar{\partial}^* u^\#\|_{s-1}^2 + \|u^\#\|_0^2 \quad s \geq 2.$$

We refer to [6] formula (1) of Main theorem as a general reference but also give an outline of the proof. For this, we have to call into play the tangential  $s$ -Sobolev norm which is defined by  $\|u\|_s = \|\Lambda^s u\|_0$ . We start from

$$(2.4) \quad \|\zeta u^\#\|_1^2 \lesssim Q(\zeta u^\#, \zeta u^\#) + \|u^\#\|_0^2;$$

this is the basic estimate for  $u^\nu$  (which vanishes at  $b\Omega$ ) whereas it is [12] Lemma 8.6 for  $u^{\tau-}$  and  $u^{\tau 0}$ . Applying (2.4) to  $\zeta' \Lambda^{s-1} \zeta u^\#$  one gets the estimate of tangential norms for any  $s$ . Finally, by non-characteristicity of  $(\bar{\partial}, \bar{\partial}^*)$  one passes from tangential to full norms along the guidelines of [16] Theorem 1.9.7. The version of this argument for  $\square$  can be found in [12] second part of p. 245. Because of (2.3), it suffices to prove (2.2) for the only  $u^{\tau+}$ . We further decompose

$$u^{\tau+} = u^{\tau+(h)} + u^{\tau+(0)},$$

where  $u^{\tau+(h)}$  is the “harmonic extension” in the sense of Kohn [12] and  $u^{\tau+(0)}$  is just the complementary part. We denote by  $\bar{\partial}^\tau$  the extension of  $\bar{\partial}_b$  from  $b\Omega$  to  $\Omega$  which stays tangential to the level surfaces  $r \equiv \text{const.}$  It acts on tangential forms  $u^\tau$  and it is defined by  $\bar{\partial}^\tau u^\tau = (\bar{\partial} u^\tau)^\tau$ . We denote by  $\bar{\partial}^{\tau*}$  its adjoint; thus  $\bar{\partial}^{\tau*} u^\tau = \bar{\partial}^*(u^\tau)$ . We use the notations  $\square^\tau$  and  $Q^\tau$  for the corresponding Laplacian and energy. We notice that over a tangential form  $u^\tau$  we have a decomposition

$$(2.5) \quad Q = Q^\tau + \|\bar{L}_n u^\tau\|_0^2.$$

The proof of (2.2) for  $u^{\tau+}$  requires two crucial technical results. Here is the first which is the most central

**Proposition 2.5.** *For the harmonic extension  $u^{\tau+(h)}$  we have*

$$(2.6) \quad \|\zeta u^{\tau+(h)}\|_s^2 \lesssim Q_{\Lambda^s \zeta'}(u^{\tau+(h)}, u^{\tau+(h)}) + \|u^{\tau+(h)}\|_0^2.$$

*Proof.* We apply compactness estimates (cf. e.g. [7] Section 6) for  $\zeta' \Lambda^s \zeta u^{\tau+(h)}$ ,

$$(2.7) \quad \|\zeta' \Lambda^s \zeta u^{\tau+(h)}\|^2 \leq \epsilon Q(\zeta' \Lambda^s \zeta u^{\tau+(h)}, \zeta' \Lambda^s \zeta u^{\tau+(h)}) + c_\epsilon \|\zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2.$$

We decompose  $Q$  according to (2.5). We calculate  $Q^\tau$  over  $\zeta' \Lambda^s \zeta u^{\tau+(h)}$  and compute errors coming from commutators  $[Q^\tau, \zeta' \Lambda^s \zeta]$ . In this calculation we assume that the cut off functions are of product type  $\zeta(z')\zeta(t)$  where  $z'$  (resp.  $t$ ) are complex (resp. totally real) tangential coordinates in  $T_{z_0} b\Omega$ . We have

$$(2.8) \quad \begin{aligned} & Q^\tau(\zeta' \Lambda^s \zeta u^{\tau+(h)}, \zeta' \Lambda^s \zeta u^{\tau+(h)}) \\ & \lesssim Q_{\zeta' \Lambda^s \zeta}^\tau(u^{\tau+(h)}, u^{\tau+(h)}) + \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta' u^{\tau+(h)}\|_{s-1}^2 \\ & + \left( \|(|\dot{\zeta}(z')| + |\dot{\zeta}'(z')|) \Lambda^s u^{\tau+(h)}\|_0^2 + \left\| \sum_{j=1}^{n-1} |r_{z_j}| (|\dot{\zeta}(t)| + |\dot{\zeta}'(t)|) \Lambda^s u^{\tau+(h)} \right\|_0^2 \right). \end{aligned}$$

We explain (2.8). First, the commutators  $[\bar{\partial}^\tau, \zeta' \Lambda^s \zeta]$  (and similarly as for  $[\bar{\partial}^{\tau*}, \zeta' \Lambda^s \zeta]$ ) are decomposed by Jacobi identity as

$$[\bar{\partial}^\tau, \zeta' \Lambda^s \zeta] = [\bar{\partial}^\tau, \zeta'] \Lambda^s \zeta + \zeta' [\bar{\partial}^\tau, \Lambda^s] \zeta + \zeta' \Lambda^s [\bar{\partial}^\tau, \zeta].$$

The central commutator  $[\bar{\partial}^\tau, \Lambda^s]$  produces the error term  $|||\zeta u^{\tau+(h)}|||_s^2$ . As for the two others, we have

$$[\bar{\partial}^\tau, \zeta(z') \zeta(t)] = [\bar{\partial}^\tau, \zeta(z')] \zeta(t) + \zeta(z') [\bar{\partial}^\tau, \zeta(t)],$$

and similarly for  $\zeta$  replaced by  $\zeta'$  and  $\bar{\partial}^\tau$  by  $\bar{\partial}^{\tau*}$ . Now,

$$(2.9) \quad [\bar{\partial}^\tau, \zeta(z')] \sim \dot{\zeta}(z').$$

On the other hand, we first notice that it is not restrictive to assume that  $\partial_{z_1}, \dots, \partial_{z_{n-1}}$  are a basis of  $T_0^{1,0} b\Omega$  for otherwise, owing to (iii), we have subelliptic estimates from which local regularity readily follows. Thus, each  $\bar{L}_j$ ,  $j = 1, \dots, n-1$ , is of type  $\bar{L}_j = r_{\bar{z}_j} \partial_{\bar{z}_n} - r_{\bar{z}_n} \partial_{\bar{z}_j}$ , and then

$$(2.10) \quad \begin{aligned} [\bar{\partial}^\tau, \zeta(t)] &\sim \sum_{j=1}^{n-1} [\bar{L}_j, \zeta(t)] \\ &\sim \sum_{j=1}^{n-1} r_{\bar{z}_j} \dot{\zeta}(t). \end{aligned}$$

By combining (2.9) with (2.10) (and using the analogous for  $\zeta'$  and  $\bar{\partial}^{\tau*}$ ), we get the last line of (2.8). This establishes (2.8). Next, since  $(\bar{\partial}^\tau, \bar{\partial}^{\tau*})$  has subelliptic estimates, say  $\eta$ -subelliptic, for  $z' \neq 0$  and hence in particular over  $\text{supp } \dot{\zeta}(z')$  and  $\text{supp } \dot{\zeta}'(z')$  and since the  $r_{\bar{z}_j}$  are, say,  $\eta$ -subelliptic multipliers even at  $z' = 0$ , then the last line of (2.8) is estimated by  $||\zeta'' \Lambda^{s-\eta} \zeta' u^{\tau+(h)}||^2$  where  $\zeta'' \equiv 1$  over  $\text{supp } \zeta'$ . This shows, using iteration over increasing  $k$  such that  $k\eta > s$  and over decreasing  $j$  from  $s-1$  to 0, that (2.7) and (2.8) imply (2.6) provided that we add on the right side the extra term  $||\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}||^2$ . Note that, as a result of the inductive process, we have to replace  $Q_{\zeta' \Lambda^s \zeta}$  in (2.8) by  $Q_{\Lambda^s \zeta'}$  in (2.6).

Up to this point the argument is the same as in [11] and does not make any use of the specific properties of the harmonic extension  $u^{\tau+(h)}$ . We start the new part which is dedicated to prove that  $||\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}||^2$  can be removed from the right of (2.6). For this we have to use the main property of this extension expressed by [12] Lemma 8.5, that is,

$$(2.11) \quad ||\bar{L}_n \zeta u^{\tau+(h)}||_0^2 \lesssim \sum_{j=1}^{n-1} ||\bar{L}_j \zeta u_b^{\tau+}||_{b, -\frac{1}{2}}^2 + ||u^{\tau+}||_0^2.$$

Note that (2.11) differs from [12] Lemma 8.5 by  $[\bar{L}_n, \Psi^+]$ ; but this is an error term which can be taken care of by  $u^{\tau+0}$  to which elliptic estimates apply. Applying (2.11) to  $\zeta' \Lambda^s \zeta u^{\tau+(h)}$  (for the first inequality below), and using the classical inequality  $\|\cdot\|_{b, -\frac{1}{2}}^2 \leq c_\epsilon \|\cdot\|_0^2 + \epsilon \|\partial_r \cdot\|_{-1}^2$  (cf. e.g. [8] (1.10)) together with the splitting  $\partial_r = \bar{L}_n + Tan$  (for the second), we get

$$\begin{aligned}
 (2.12) \quad \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 &\stackrel{\text{by (2.11)}}{\lesssim} \sum_{j=1}^{n-1} \|\bar{L}_j \zeta' \Lambda^s \zeta u_b^{\tau+}\|_{b, -\frac{1}{2}}^2 + \|\zeta' \Lambda^s \zeta u^{\tau+}\|_0^2 \\
 &\lesssim c_\epsilon \sum_{j=1}^{n-1} \|\bar{L}_j \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 + \epsilon \sum_{j=1}^{n-1} \|[\bar{L}_n, \bar{L}_j] \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 \\
 &\quad + \epsilon \sum_{j=1}^{n-1} \|Tan \bar{L}_j \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 + \|\zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2.
 \end{aligned}$$

The first term on the right of the last inequality is controlled by  $\sum_{j=1}^{n-1} \|\zeta' \Lambda^s \zeta \bar{L}_j u^{\tau+(h)}\|^2 + \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta'' u^{\tau+(h)}\|_{s-1}^2$  by the first part of the proposition; moreover, we have the immediate estimate  $\sum_{j=1}^{n-1} \|\zeta' \Lambda^s \zeta \bar{L}_j u^{\tau+(h)}\|^2 \lesssim Q_{\Lambda^s \zeta'}^\tau(u^{\tau+(h)}, u^{\tau+(h)})$ . The term which carries  $\epsilon Tan$ , after  $Tan$  has been annihilated by the Sobolev norm of index  $-1$ , has the same estimate as the first term. It remains to control the second term in the right which involves  $\epsilon \bar{L}_n$ . We rewrite  $\bar{L}_n \bar{L}_j = \bar{L}_j \bar{L}_n + [\bar{L}_n, \bar{L}_j]$ ; when  $\bar{L}_j$  moves in first position, it is annihilated by  $-1$  and what remains is absorbed in the left. As for the commutator, we have

$$\begin{aligned}
 \|[\bar{L}_n, \bar{L}_j] \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 &\lesssim \|\zeta u^{\tau+(h)}\|_s^2 + \|\partial_r \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2 \\
 &\lesssim \|\zeta u^{\tau+(h)}\|_s^2 + \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_{-1}^2,
 \end{aligned}$$

where we have used the splitting  $\partial_r = Tan + \bar{L}_n$  in the second inequality. Again, the term with  $\bar{L}_n$ , which now comes in  $-1$  norm, is absorbed in the left of (2.12). Summarizing up, we have got

$$\begin{aligned}
 (2.13) \quad \|\bar{L}_n \zeta' \Lambda^s \zeta u^{\tau+(h)}\|_0^2 &\lesssim c_\epsilon Q_{\Lambda^s \zeta'}^\tau(u^{\tau+(h)}, u^{\tau+(h)}) \\
 &\quad + \|\zeta u^{\tau+(h)}\|_s^2 + \|\zeta'' u^{\tau+(h)}\|_{s-1}^2.
 \end{aligned}$$

But  $\|[\bar{L}_n, \cdot]\|^2$  comes with a factor  $\epsilon$  of compactness and hence the term in  $s$ -norm in the last line can be absorbed in the left of the initial

inequalities (2.7) or (2.6). Finally, we use an inductive argument to go down from  $s - 1$  to 0. This concludes the proof of the proposition.  $\square$

We remark now that

$$\begin{aligned}
 (2.14) \quad & ||\zeta u^{\tau+(h)}||_0^2 \lesssim ||\zeta u_b^{\tau+}||_{b, -\frac{1}{2}}^2 \\
 & \lesssim ||\zeta u^{\tau+}||_0^2 + |||\partial_r \zeta u^{\tau+}|||_{-1}^2 \\
 & \lesssim ||\zeta u^{\tau+}||_0^2 + |||\bar{L}_n \zeta u^{\tau+}|||_{-1}^2 + |||Tan \zeta u^{\tau+}|||_{-1}^2 \\
 & \lesssim Q_{\Lambda^{-1}\zeta}(u^{\tau+}, u^{\tau+}) + ||\zeta u^{\tau+}||_0^2.
 \end{aligned}$$

The same inequality also holds for  $u^{\tau+(h)}$  replaced by  $u^{\tau+(0)}$  on account of the identity  $u^{\tau+(0)} = u^{\tau+} + u^{\tau+(h)}$ . We need another preparation result

**Proposition 2.6.** *We have*

$$(2.15) \quad Q_{\Lambda^s \zeta'}(u^{\tau+(h)}, u^{\tau+(h)}) \lesssim Q_{\Lambda^s \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+})$$

and

$$\begin{aligned}
 (2.16) \quad & |||\zeta u^{\tau+(0)}|||_s^2 \lesssim Q_{\Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-2} \zeta'}(u^{\tau+}, u^{\tau+}) \\
 & + |||\zeta' \Delta u^{\tau+}|||_{s-2}^2 + ||u^{\tau+}||_0^2.
 \end{aligned}$$

*Proof.* The proof of (2.15) is an immediate combination of the formulas  $||\zeta' u^{\tau+(h)}||_0 \lesssim ||\zeta' u_b^{\tau+}||_{b, -\frac{1}{2}}$  and  $||\zeta' u^{\tau+}||_{b, -\frac{1}{2}} \lesssim ||\zeta' u^{\tau+}||_0 + |||\partial_r \zeta' u^{\tau+}|||_{-1}^2$ .

We prove now (2.16). By elliptic estimate for  $u^{\tau+(0)}$  (which vanishes at  $b\Omega$ ) with respect to the order 2 elliptic operator  $\Delta$ , we have

$$(2.17) \quad |||\zeta u^{\tau+(0)}|||_s^2 \lesssim |||\zeta' \Delta u^{\tau+(0)}|||_{s-2}^2 + ||u^{\tau+(0)}||_0^2.$$

This result of Sobolev regularity at the boundary is very classical: it is formulated, for functions in  $H_0^1$  such as the coefficients of  $u^{\tau+(0)}$ , e.g. in Evans [5] Theorem 5 p. 323. Owing to the identity  $\Delta u^{\tau+(0)} = \Delta u^{\tau+} + P^1 u^{\tau+(h)}$  for a 1-order operator  $P^1$  (cf. [12] p. 241), we can replace  $\Delta u^{\tau+(0)}$  by  $\Delta u^{\tau+}$  on the right side of (2.17) putting the contribution of  $P^1$  into an error term of type  $|||\zeta' u^{\tau+(h)}|||_{s-1} + |||\zeta' \partial_r u^{\tau+(h)}|||_{s-2}$ , which can be estimated, on account of the splitting  $\partial_r = \bar{L}_n + Tan$ , by  $|||\zeta' u^{\tau+(h)}|||_{s-1} + |||\zeta'' u^{\tau+(h)}|||_{s-2} + Q_{\Lambda^{s-2} \zeta'}(u^{\tau+(h)}, u^{\tau+(h)})$ . We write the terms of order  $s - 1$  and  $s - 2$  as a common  $|||\zeta'' u^{\tau+(h)}|||_{s-1}$  that we can estimate, using (2.6) and (2.15), by

$$|||\zeta'' u^{\tau+(h)}|||_{s-1}^2 \lesssim Q_{\Lambda^{s-1} \zeta'''}(u^{\tau+}, u^{\tau+}) + Q_{\Lambda^{s-2} \partial_r \zeta'''}(u^{\tau+}, u^{\tau+}).$$



This brings down from  $s - 1$  to 0 the Sobolev index in the error term. This 0-order term  $\|u^{\tau+(h)}\|_0^2$ , together with its companion  $\|u^{\tau+(0)}\|_0^2$  in the right of (2.17), is estimated, because of (2.14), by  $\|u^{\tau+}\|_0^2$  up to a term  $Q_{\Lambda^{-1}\zeta}$  which is controlled by the right side of (2.16). This concludes the proof of (2.16).  $\square$

*End of proof of Theorem 2.3.* We prove (2.2) for  $u^{\tau+}$ ; this implies the conclusion in full generality according to the first part of the proof. We have

$$\begin{aligned}
 (2.18) \quad |||\zeta u^{\tau+(h)}|||_s^2 &\underset{\sim}{\leq} Q_{\Lambda^s \zeta'}(u^{\tau+(h)}, u^{\tau+(h)}) + \|u^{\tau+(h)}\|_0^2 \\
 &\underset{\sim}{\leq} Q_{\Lambda^s \zeta'}(u^{\tau+}, u^{\tau+}) + Q_{\partial_r \Lambda^{s-1} \zeta'}(u^{\tau+}, u^{\tau+}) + \|u^{\tau+}\|_0^2. \\
 &\quad \text{by (2.6)} \\
 &\quad \text{by (2.15) and (2.14)}
 \end{aligned}$$

We combine (2.18) with (2.16); what we get is

$$\begin{aligned}
 (2.19) \quad |||\zeta u^{\tau+}|||_s^2 &\leq |||\zeta u^{\tau+(h)}|||_s^2 + |||\zeta u^{\tau+(0)}|||_s^2 \\
 &\underset{\sim}{\leq} \|\zeta' \bar{\partial} u^{\tau+}\|_s^2 + \|\zeta' \bar{\partial}^* u^{\tau+}\|_s^2 + |||\zeta' \Delta u^{\tau+}|||_{s-2}^2 + \|u^{\tau+}\|_0^2.
 \end{aligned}$$

By the non-characteristicity of  $Q$ , we can replace the tangential norm  $|||\cdot|||_s$  by the full norm  $\|\cdot\|_s$  in the left of (2.19). (The explanation of this point can be found, for example, in [12] second part of p. 245.) This proves (2.2) for  $u^{\tau+}$  and thus also for a general  $u$ .  $\square$

We modify  $b\Omega$  outside a neighborhood of  $z_o$  where it satisfies the hypotheses of Theorem 1.1 so that it is strongly pseudoconvex in the modified portion and bounds a relatively compact domain; in particular, there is well defined the  $H^0$  inverse  $N$  of  $\square$  in this domain. There is an immediate crucial consequence of Theorem 2.3.

**Theorem 2.7.** *We have that*

$$(2.20) \quad \bar{\partial}^* N \text{ is exactly regular over } \text{Ker } \bar{\partial}$$

and

$$(2.21) \quad \bar{\partial} N \text{ is exactly regular over } \text{Ker } \bar{\partial}^*.$$

*Proof.* As for (2.20), we put  $u = \bar{\partial}^* N f$  for  $f \in \text{Ker } \bar{\partial}$ . We get

$$\begin{cases} \bar{\partial} u = f, \\ \bar{\partial}^* u = 0, \\ \Delta u = (\vartheta \bar{\partial} + \bar{\partial} \vartheta) \bar{\partial}^* N f \\ \quad = \vartheta (\bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial}) N f + \bar{\partial} \vartheta \bar{\partial}^* N f \\ \quad = \vartheta \square N f = \vartheta f. \end{cases}$$

Thus, by (2.2)

$$\begin{aligned} (2.22) \quad \|\zeta u\|_s^2 &\lesssim \|\zeta' f\|_s^2 + \|\zeta' \vartheta f\|_{s-2}^2 + \|u\|_0^2 \\ &\lesssim \|\zeta' f\|_s^2 + \|u\|_0^2. \end{aligned}$$

To prove (2.21), we put  $u = \bar{\partial} N f$  for  $f \in \text{Ker } \bar{\partial}^*$ . We have a similar calculation as above which leads to the same formula as (2.22) (with the only difference that  $\vartheta$  is replaced by  $\bar{\partial}$  in the intermediate inequality). Thus from (2.22) applied both for  $\bar{\partial}^* N$  and  $\bar{\partial} N$  on  $\text{Ker } \bar{\partial}$  and  $\text{Ker } \bar{\partial}^*$  respectively, we conclude that these operators are exactly regular.  $\square$

We are ready for the proof of Theorem 1.1. This follows from Theorem 2.7 by the method of Boas-Straube.

*Proof of Theorem 1.1.* From the regularity of  $\bar{\partial}^* N$  it follows that the Bergman projection  $B$  is also regular. (Notice that exact regularity is perhaps lost by taking  $\bar{\partial}$  in  $B$ .) We exploit formula (5.36) in [15] in unweighted norms, that is, for  $t = 0$ :

$$\begin{aligned} N_q &= B_q(N_q \bar{\partial})(\text{Id} - B_{q-1})(\bar{\partial}^* N_q) B_q \\ &\quad + (\text{Id} - B_q)(\bar{\partial}^* N_{q+1}) B_{q+1}(N_{q+1} \bar{\partial})(\text{Id} - B_q). \end{aligned}$$

Now, in the right side, the  $\bar{\partial} N$ 's and  $\bar{\partial}^* N$ 's are evaluated over  $\text{Ker } \bar{\partial}^*$  and  $\text{Ker } \bar{\partial}$  respectively; thus they are exactly regular. The  $B$ 's are also regular and therefore such is  $N$ . This concludes the proof of Theorem 1.1.  $\square$

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